

- Investigated quantities:
 - Rotational angular velocity
 - Precessional rotational velocity
 - Angle of precession
 - Trajectory of coin

ω_3 – rotational angular velocity

ω_2 – precessional angular velocity

φ – angle of precession

x, y, z – axes of the coin coordinate system

Used coins

mass [g]	radius [cm]	thickness [cm]
7.3	1.22	0.19

mass [g]	radius [cm]	thickness [cm]
3.2	1	0.15

mass [g]	radius [cm]	thickness [cm]
67.3	4.75	0.70

mass [g]	radius [cm]	thickness [cm]
102.5	4.9	2



Ways of starting the coin motion

- Two ways of starting the coin motion exist:
 1. No initial horizontal momentum



- The motion is started from above

2. With a significant initial horizontal momentum



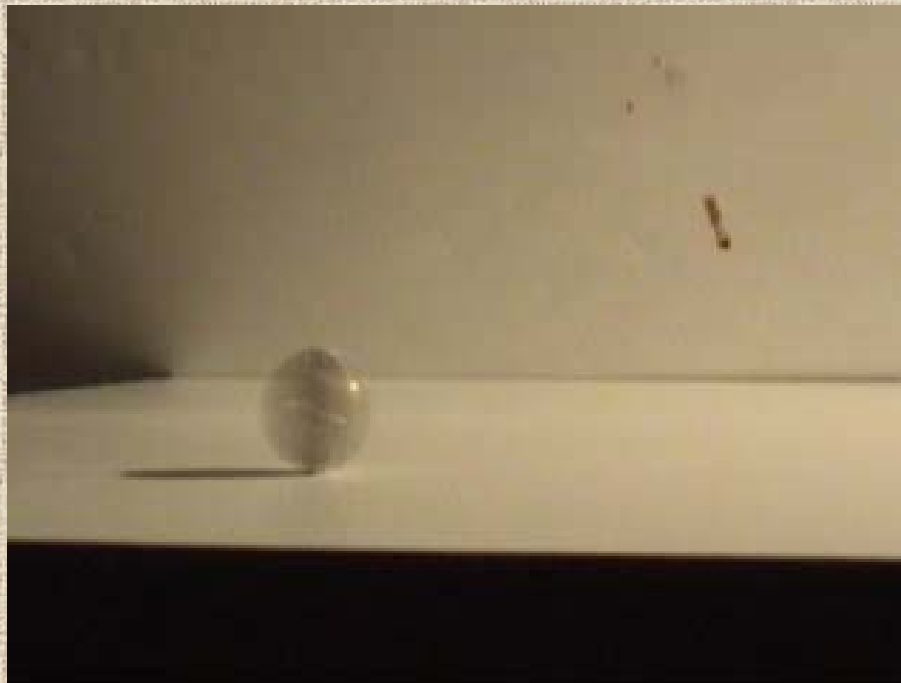
- The motion is started with a side kick

1. The coin has no horizontal velocity

Phases of motion:

1. Slow precession

- the angle of precession changes slowly with time



2. Slipping

- The precessional angle drops abruptly
- Cause: slipping (the gyroscopic effect and friction cannot hold the coin straight any more)



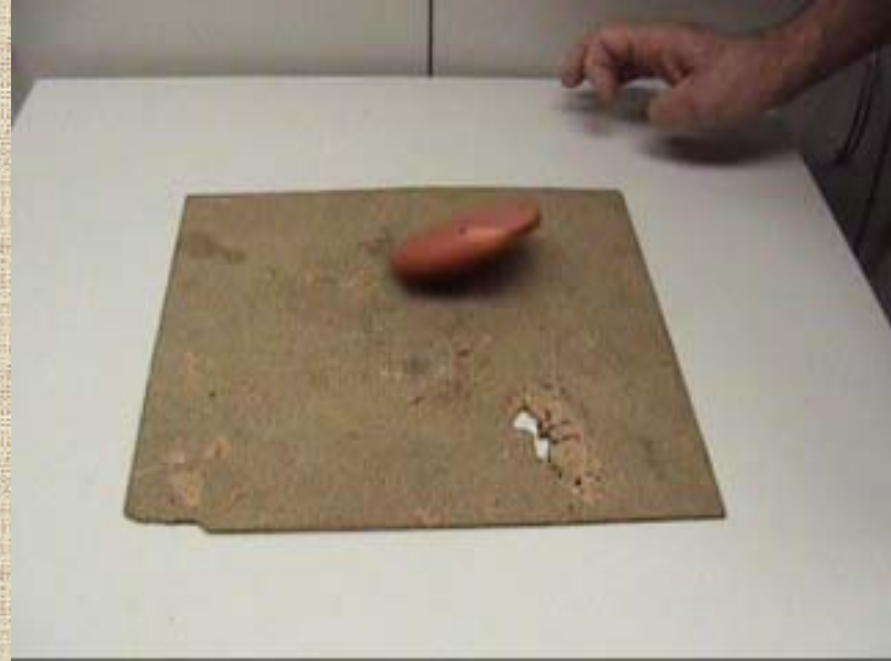
3. Dancing

- The most complex phase
- The coin doesn't spin any more
- Because of energy losses the edges start to hit the surface
- Thus the characteristic sound is obtained
- The duration of this phase can be quite long depending on the elasticity of the surface and coin:

Dancing on
plastic/wood – elastic
collisions



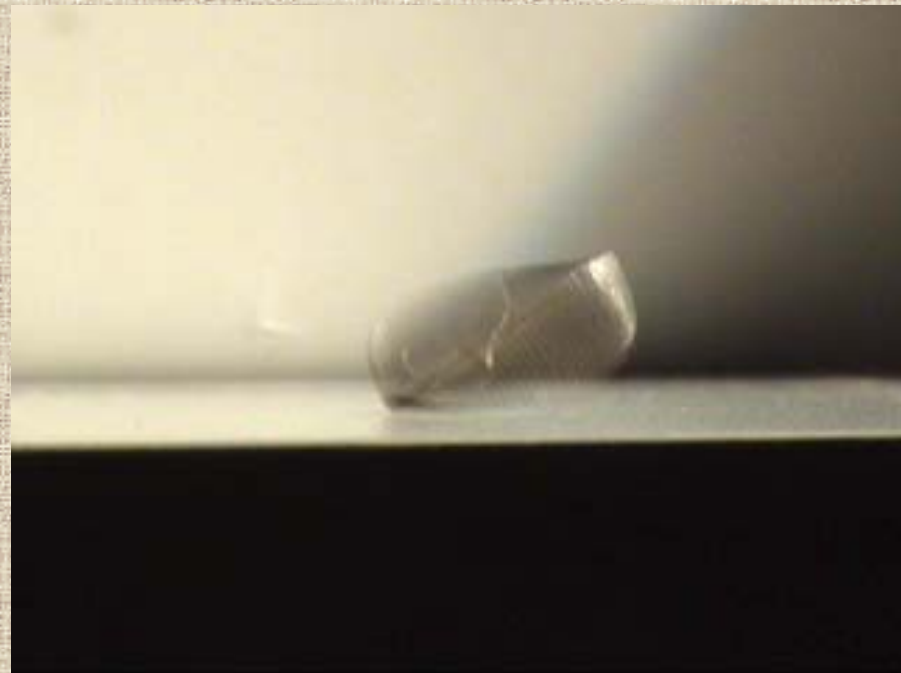
Dancing on cork –
great energy losses
in collisions



Dancing on cardboard
– semielastic collisions



Dancing of elastic
metal coin



2. The coin is given an initial horizontal velocity

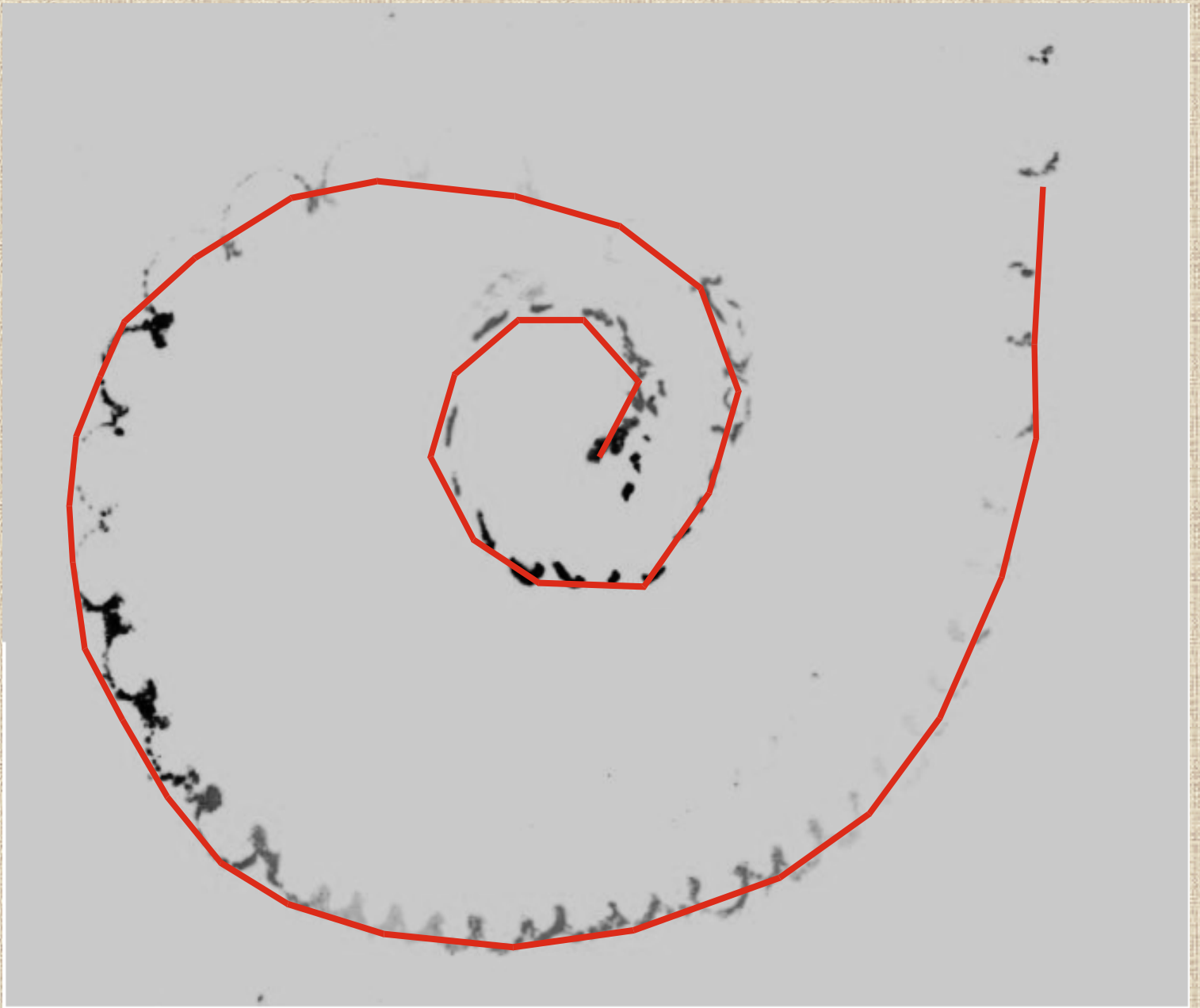
- **Phases of motion**

- Neglecting the translatory motion, the phases are the same as in the previous case



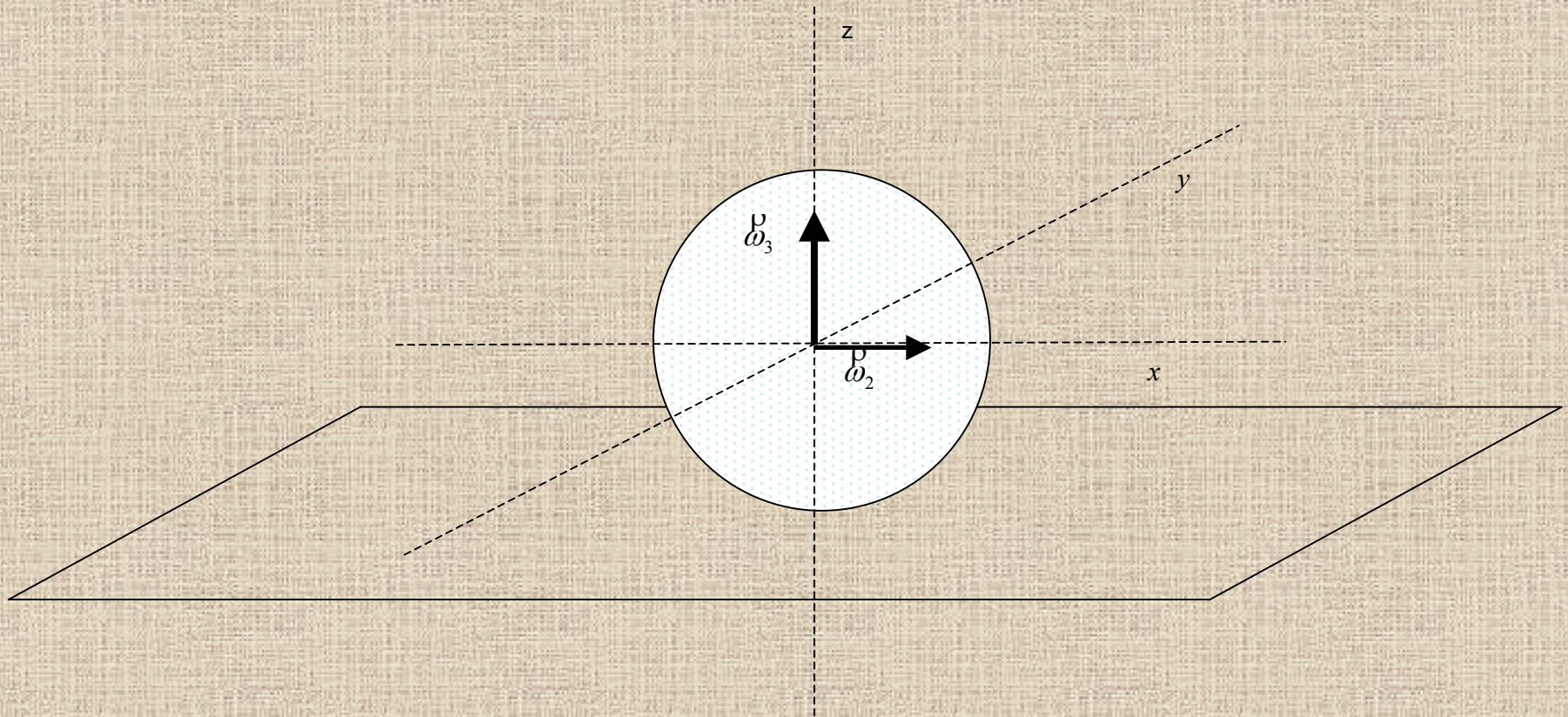
Trajectory

- Because of the precession the initially linear trajectory gets curved
- The length of the linear part depends on the initial momentum of the coin
- The best curve fit for the whole trajectory is a logarithmic spiral
- The trajectory graph was obtained by letting the coin move on a sooty plate and filming the trace it left:



Model of coin motion

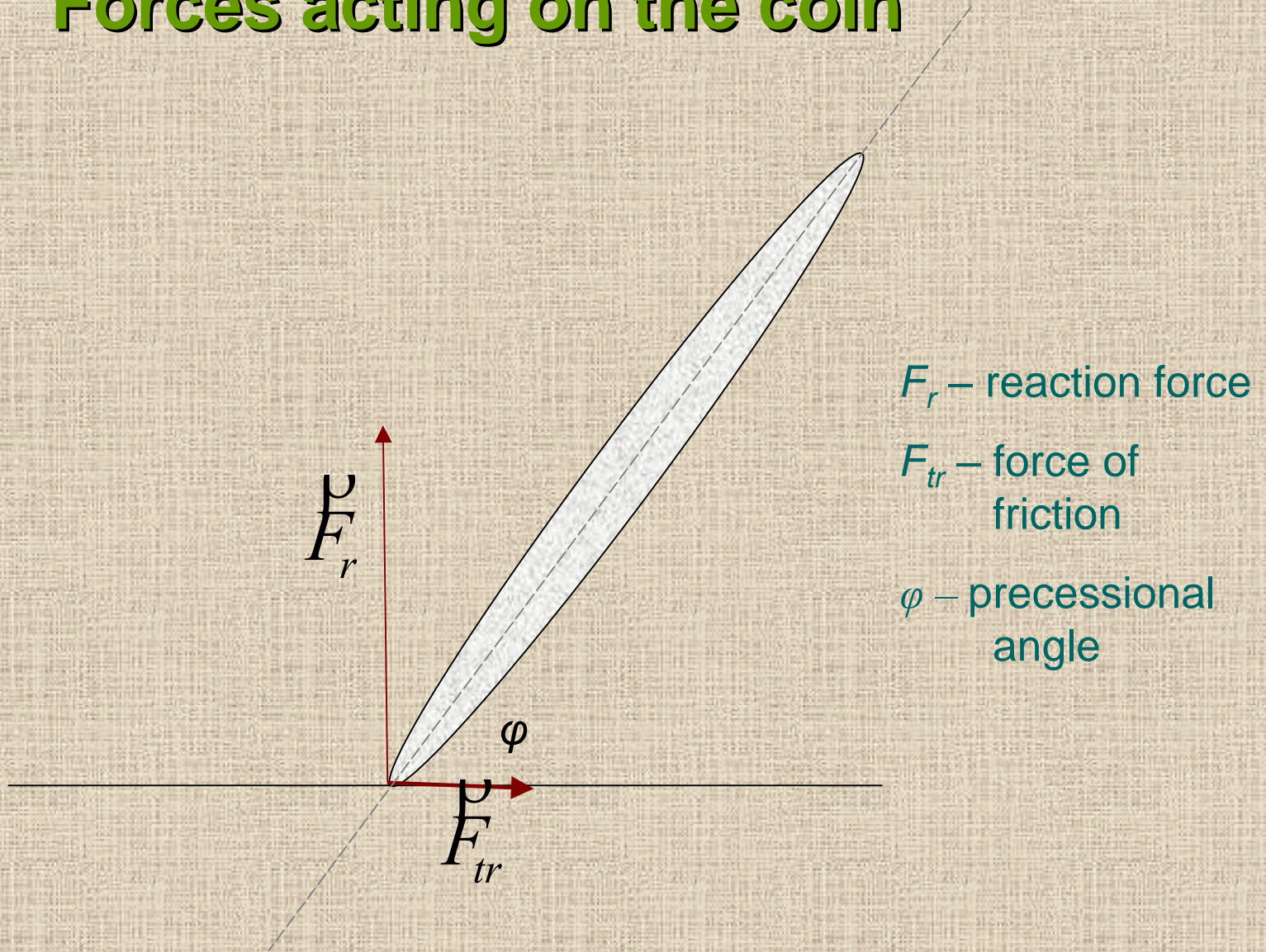
Coin coordinate system



ω_3 – rotational angular velocity

ω_2 – precessional angular velocity

Forces acting on the coin



Euler equations

$$I_1 \frac{d\omega_1}{dt} - (I_2 - I_3)\omega_2\omega_3 = \tau_1$$

$$I_2 \frac{d\omega_2}{dt} - (I_3 - I_1)\omega_3\omega_1 = \tau_2$$

$$I_3 \frac{d\omega_3}{dt} - (I_1 - I_2)\omega_1\omega_2 = \tau_3$$

ω_i – angular velocity components

τ_i – external torque components

I_i – principal moments of inertia

- Neglecting friction the torque components become

$$\tau_1 = \tau_3 = 0$$

$$\tau_2 = \tau_r \quad \tau_r \text{ – reaction torque}$$

Euler equations cont.

$$I_1 \frac{d\omega_1}{dt} = 0$$

$$\Rightarrow I_2 \frac{d\omega_2}{dt} - (I_2 - I_1)\omega_3\omega_1 = \tau_r$$

$$I_2 \frac{d\omega_3}{dt} - (I_1 - I_2)\omega_1\omega_2 = 0$$

- Initial conditions: $\omega_1(t) = 0$

$$\omega_3 = \omega_0$$

- Reaction torque: $|\tau_r| = -mgr \cos \varphi$

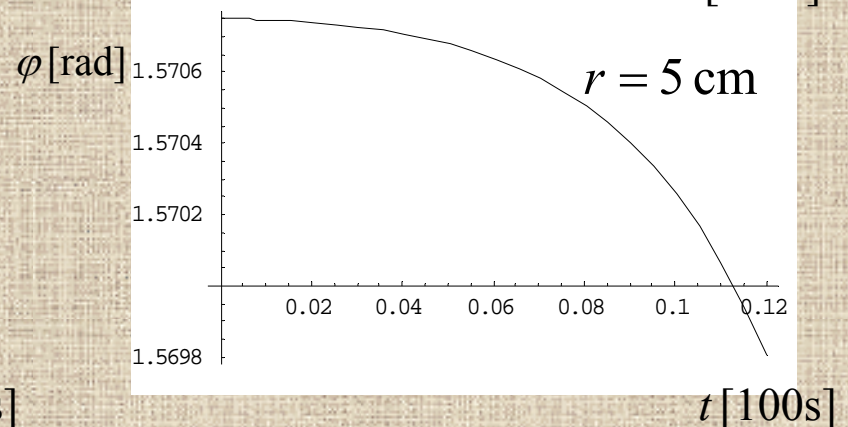
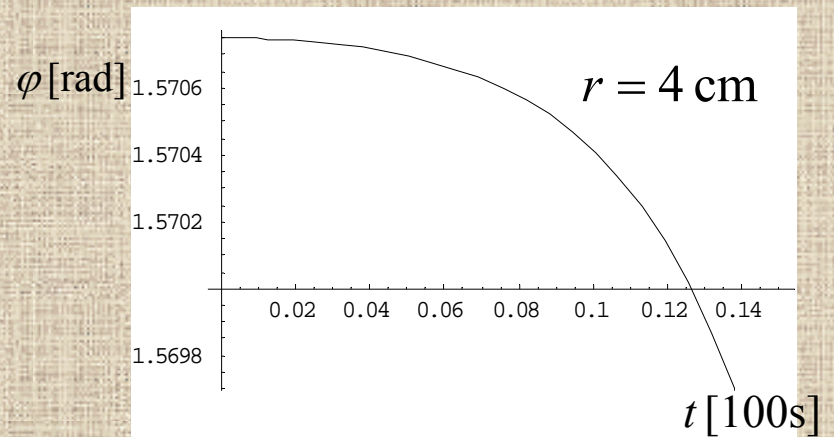
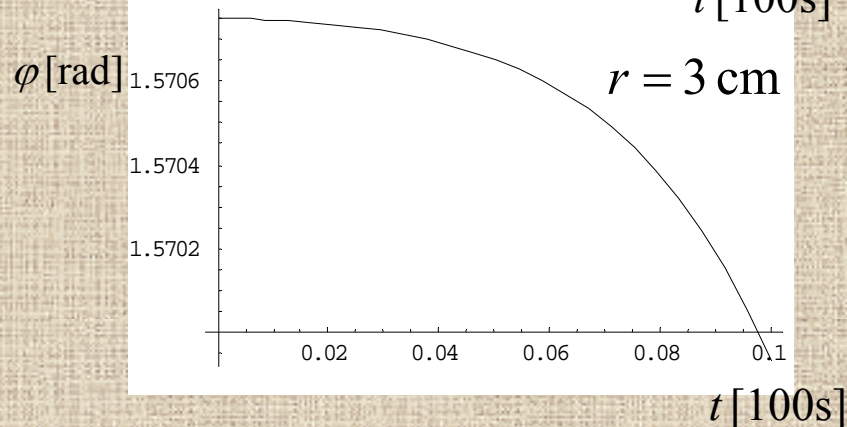
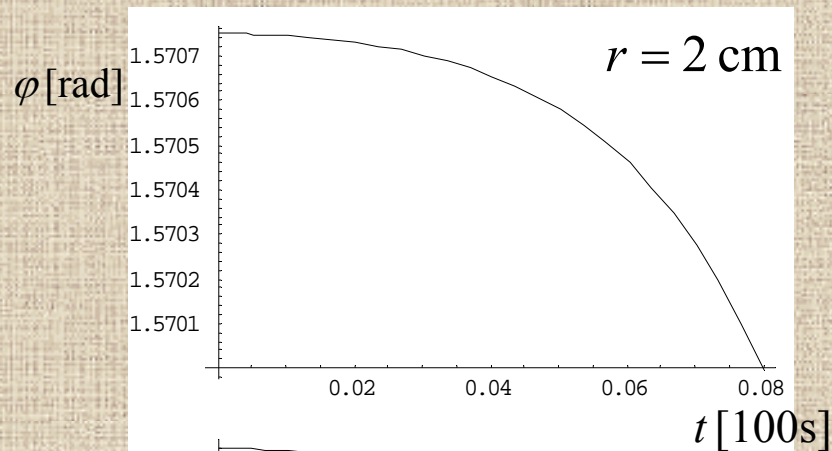
$$\Rightarrow \frac{d^2\varphi}{dt^2} = -\frac{4g}{r} \cos \varphi$$

φ – precessional angle

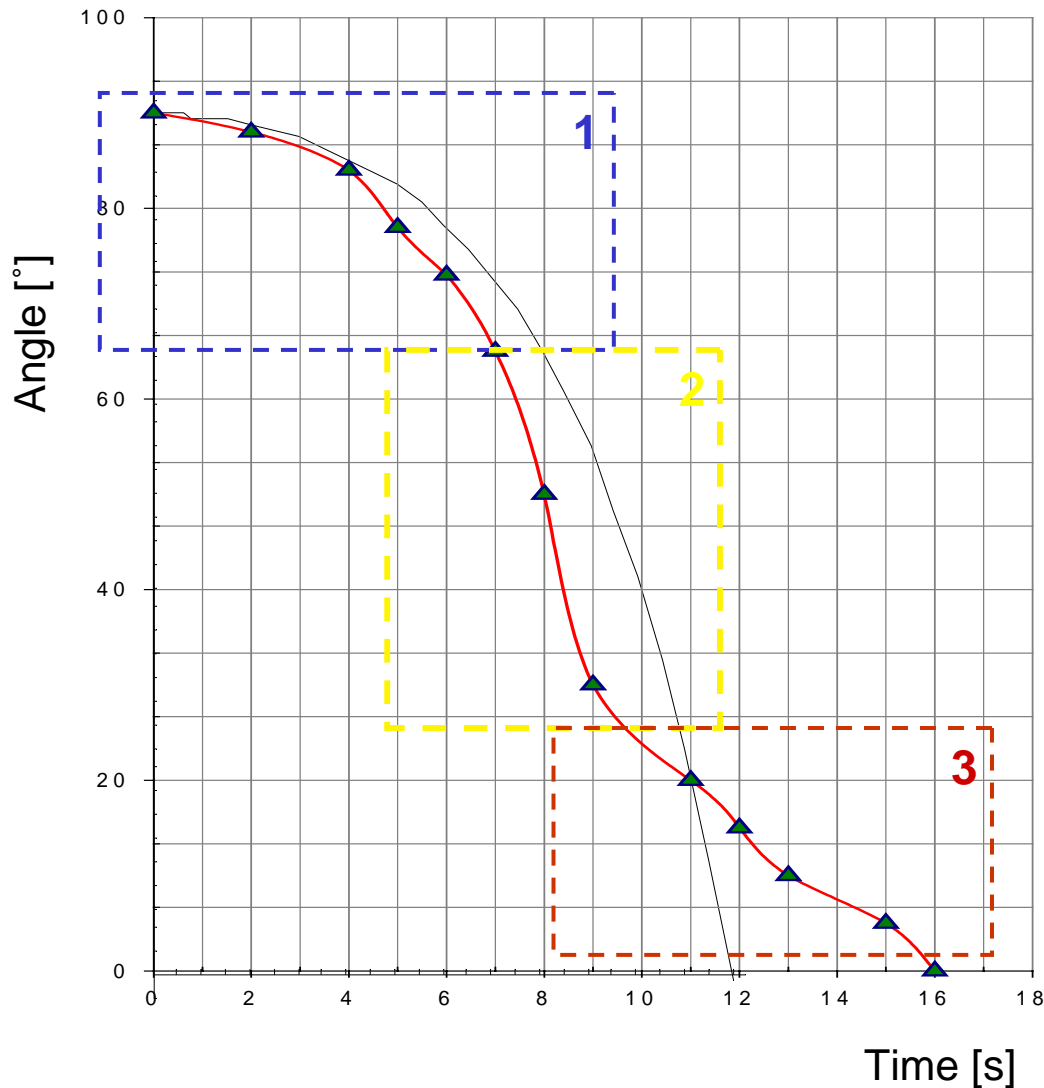
r – coin radius

Solution of Euler equation

- The precessional angle equation isn't soluble in closed form
- Therefore numeric integration was performed:



Fitting the theoretical curve



1. Agreement
2. Slipping
(error ~20%)
3. "dancing" of the coin
(feeble agreement)

Conclusion

- Two ways of coin motion were investigated
- A mathematical model of motion was proposed
- The motion can be divided in three parts
 1. Large precessional angle – agreement with theory is very good (error $\sim 5\%$)
 2. Slipping – the agreement is not as well because of the coin slipping
 3. “dancing” - agreement is not good because of the additional effects (collisions with the floor)

Add – on: Initial destabilizing torque

